# ON ROTATIONAL SPRING CONSTANTS AT THE JUNCTURE OF A RADIAL NOZZLE AND A SPHERICAL SHELL 

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## ABSTRACT

This paper presents the values of the rotational spring constants (i.e. Spring Constants due to bending moment) at the juncture of a radial nozzle and a spherical shell. This rotational spring constant is a function of diameters and thicknesses of nozzle and vessel. It is known that a lower Rotational Spring Constant at the juncture would reduce the local peak stresses which is very important in fatigue design of pressure vessel. This study shows that the value of the Rotational Spring Constant may be reduced by increasing the shell diameter and decreasing the shell thickness, nozzle diameter and nozzle thickness. Analytical derivations for this spring constant are presented in this paper. The actual values of Rotational Spring Constants for various nozzle-shell combination have been computed using a digital computer. These values are tabulated and plotted in this paper to facilitate the vessel design and stress analysis.

## NOMENCLATURE


$K_{B}=$ Rotational Spring Constant
$\mathrm{K}_{1}$ to $\mathrm{K}_{5}=$ Constants
Ker $s, K e i s, K e r u, K e i u=$
Kelvin functions of zero order
Ker's, Kei's, Ker'u, Kei'u =
Derivatives of Kelvin functions
$1=\left[R^{2} t^{2} /\left(12\left(1-\mu^{2}\right)\right)\right]^{\frac{1}{4}}$
$\mathrm{L}_{1}$ to $\mathrm{L}_{4}=$ Constants
$M=$ Bending moment coming through a nozzle
$M_{0}=$ Bending moment $\left(M_{x}\right)$ in the nozzle at $x=0$
$M_{x}=$ Radial moment acting per unit width upon a normal section of the spherical shell
$M_{y}=$ Tangential moment acting per unit width upon a meridional section of the spherical shell
$M_{x y}=$ Twisting moment in the nozzle or spherical shell
$\mathrm{N}=\mathrm{Eh}^{3} /\left[12\left(1-\mu^{2}\right)\right.$ Flexural rigidity of nozzle
$\mathrm{N}_{\mathrm{x}}=$ Radial membrane force, acting per unit width upon a normal section of the spherical shell
$\mathrm{N}_{\mathrm{y}}=$ Tangential membrane force, acting per unit width upon a meridional section of the spherical shell
$Q_{x}=$ Transverse shear force in cross-section of the nozzle
$\bar{Q}_{x}=$ Equivalent transverse shear force for nozzle in section upon which $Q_{x}$ acts
$Q_{x y}=$ Transverse shear force in cross-section of the spherical shell
$\bar{Q}_{x v}=$ Equivalent transverse shear force for spherical shell in section upon which $Q_{x v}$ acts
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$R=$ Radius of middle plane of spherical shell
$r=$ Radius of curvature of middle surface of Spherical Shell in a latitudinal plane (Refer Fig. 1)
$s=r / 1=1.81784(r / R) \cdot(R / t)^{1 / 2}$
$\mathrm{t}=$ Thickness of the spherical shell
$T=$ Temperature of operation
$T_{x}=$ Axial membrane force in the nozzle
$T_{y}=$ Tangential membrane force in the nozzle
$\mathrm{T}_{\mathrm{x} \theta}=$ Membrane shear force in the nozzle
$u=a / 1=1.81784(a / R) \cdot(R / t)^{\frac{1}{2}}$
$u_{1}=$ Axial displacement of the nozzle
$v_{1}=$ Displacement of the nozzle in the tangential. ( $\theta$ ) direction
$v_{0}=$ Transverse shear $v_{x}$ in nozzle at $x=0$
$w_{1}=$ Radial deflection of the nozzle, positive if directed outwards
$w_{v}=$ Radial deflection of the spherical shell, positive if directed outwards
$x=$ Axial coordinate of cylindrical shell (nozzle)
$y=$ Tangential coordinate of cylindrical and spherical shells
$z=$ Radial coordinate of cylindrical shell (nozzle)

## Greek Symbols

$\alpha=$ Slope angle at the juncture of spherical shell and nozzle with respect to their original positions due to their deflections
$\alpha_{1}=(1 / a)\left[1-(1 / 2 \mu)+\left(3\left(1-\mu^{2}\right) \boldsymbol{\varkappa}^{2}+1-\left(3 / 4 \mu^{2}\right)\right)^{\frac{1}{2}}\right]^{\frac{1}{2}}$
$\beta=R / t$ (A Shell Parameter)
$\beta_{0}=\left[3\left(1-\mu^{2}\right) /\left(a^{2} h^{2}\right)\right]^{\frac{1}{4}}$
$\beta_{1}=(1 / a)\left[-1+(1 / 2 \mu)+\left(3\left(1-\mu^{2}\right) r^{2}+1-\left(3 / 4 \mu^{2}\right)\right)^{\frac{1}{2}}\right)^{\frac{1}{2}}$
$=a / h$ (A Nozzle Parameter)
$=$ Unit strain
= Polar coordinate for spherical shell and nozzle, in radians
$=$ Poisson's Ratio (Assumed as 0.3 for calculations)
$\rho=t / h$ (Shell Thickness/Nozzle Thickness)
$\phi=$ Angle between shell axis and normal to middle surface of spherical shell, in radians
$\nabla^{2}=\left(\mathrm{d}^{2} / \mathrm{dr}^{2}\right)+(1 / \mathrm{r})(\mathrm{d} / \mathrm{dr})$

## INTRODUCTION

With the advent of more and more chemical plants, power plants and other process plants, the safety requirement and economy in design of pressure vessels are gaining high priority. This leads to the requirement of a more accurate stress analysis.

It is known that there exist highly localized stresses at the juncture of nozzle and pressure vessels, both in cylindrical and spherical tupes. However, these stresses cannot be accurately evaluated even with most modern and highly sophisticated computer programs without reliable values of the spring constants at these junctures. The radial spring constants at the juncture of a radial nozzle and a spherical shell are given in Ref [1]. This paper deals specifically with rotational spring constants (i.e. spring constants due to bending moment) at the juncture of a nozzle and a spherical shell.

Bijlaard has studied the differential equations for a bending moment acting on a spherical shell. His solution leads to the deflection of the spherical shell only [2]. He did not specify the relationships of rotational spring constants in terms of shell parameters (diameter and thickness) and the nozzle parameters (diameter and thickness). This part is explicitly developed here. The solution involves Kelvin functions of zero and first order along with their derivatives. These are obtained by programming the mathematical equations given by Abromowitz and Stegun [3].

DERIVATION OF EXPRESSION FOR DEFLECTION OF SPHERICAL SHELL DUE TO BENDING MOMENT

The radial deflection and stress function for the case of bending moment acting on the nozzle is given by [4]:

$$
\begin{align*}
& { }^{w}{ }_{v}=\left(C_{3} \text { Ker's }^{\prime}+C_{4} \text { Kei's }^{\prime}\right) \operatorname{Cos} \theta  \tag{1}\\
& \mathrm{F}=\underset{\mathrm{Cos} \theta}{ } \quad\left[\mathrm{Et}^{2} /\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}}\right] \cdot\left(\mathrm{C}_{3} \mathrm{Kei}^{\prime} \mathrm{s}-\mathrm{C}_{4} \text { Ker }^{\prime} \mathrm{s}+\mathrm{C}_{12} \mathrm{~s}^{-1}\right)  \tag{2}\\
& \text { where } w_{v}=\text { Radial deflection of spherical shell } \\
& \text { (in the direction of exterior normal) } \\
& \mathrm{F}=\text { Stress function } \\
& \theta=\text { Polar co-ordinate for cylindrical and spherical } \\
& \text { shells, in radians }
\end{align*}
$$



Fig. 1 Spherical shell subjected to a bending monent (M) acting on a nozzle

The deflections of the nozzle in axial, tangential and radial directions are given by [2]:

$$
\begin{align*}
u_{1}= & {\left[e^{-\alpha 1 x}\left(H_{5} \operatorname{Cos} \beta_{1} x+H_{6} \operatorname{Sin} \beta_{1} x\right)-\right.} \\
& \left.(M /(\pi a h E)) \cdot(x / a)-H_{3} a\right] \cdot \operatorname{Cos} \theta  \tag{3}\\
v_{1}= & {\left[e^{-\alpha_{1} x}\left(H_{7} \operatorname{Cos} \beta_{1} x+H_{8} \operatorname{Sin} \beta_{1} x\right)-\right.} \\
& \left.(M /(2 \pi a h E)) \cdot\left(x^{2} / a^{2}\right)-H_{3} x-H_{4}\right]  \tag{4}\\
w_{1}= & {\left[e^{-\alpha_{1} x}\left(H_{1} \operatorname{Cos} \beta_{1} x+H_{2} \operatorname{Sin} \beta_{1} x\right)+\right.} \\
& \left.\left.(M /(2 \pi a h E)) \cdot\left(x^{2} / a^{2}\right)+2 \mu\right)+H_{3} x+H_{4}\right] \tag{5}
\end{align*}
$$

Equations (1) and (2) contain 3 unknown constants, namely $C_{3}, C_{4}$ and $C_{12}$ while equations (3), (4) and (5) contain 4 unknown constants, namely $\mathrm{H}_{1}, \mathrm{H}_{2}, \mathrm{H}_{3}$ and $\mathrm{H}_{4}$. (The constants $\mathrm{H}_{5}, \mathrm{H}_{6}, \mathrm{H}_{7}$ and $\mathrm{H}_{8}$ can be expressed in terms of constants $\mathrm{H}_{1}$ and $\mathrm{H}_{2}$ as per equations (27) and (40) of Ref. [2]). These 7 unknown constants require 7 boundary conditions between the spherical shell and the nozzle. These are given as follows:

Boundary Condition I:
The cross-section of the nozzle remains circular and its radius does not change. Hence B.C. I is given by:
$\left(\mathrm{w}_{1}\right)_{\mathrm{x}=0}=0$
Boundary Condition II:
Since the plane cross-section remains the same, the B.C. II is given by:
$\left(u_{1}\right)_{x=0}=0$


Fig. 2 Deflections of nozzle and spherical shell due to a bending moment


Fig. 3 Normal and shear forces acting on the spherical. shell at the nozzle-sphrical shell juncture

Boundary Condition III:
Since the shell is subjected to the external moment, we have from Ref. [5]:
$c_{12}=\left[3\left(1-\mu^{2}\right)\right]^{\frac{1}{2}} \cdot \mathrm{RM} /\left(\pi E t^{2} 1\right)$
Boundary Condition IV:
From Fig. 2, the compatibility of rotation at the juncture of spherical shell and nozzle requires:
$-\left(\partial \mathrm{w}_{\mathrm{v}} / \partial \mathrm{r}\right)-\left(\partial \mathrm{w}_{1} / \partial \mathrm{x}\right)=-\mathrm{w}_{\mathrm{v}} / \mathrm{a}$
or, $\left(w_{v} / a-\partial w_{v} / \partial r\right)_{r=a}=\left(\partial w_{1} / \partial x\right)_{x=0}$
Boundary Condition V:
The rotation of the spherical shell at $r=a$ and that of nozzle at $x=0$ should be equal. This leads to:
$\left(\varepsilon_{y v}\right)_{r=a}=\left(\varepsilon_{\theta}\right)_{x=0}$
Boundary Condition VI:
At the juncture, the bending moment $M_{x}$ in the walls of the nozzle and shell should be equal. Hence, we get:
$\left(M_{X V}\right)_{r=a}=\left(M_{x}\right)_{x=0}$
Boundary Condition VII:
The shear force $Q_{X}$ in the nozzle at $x=0$ has to be in equilibrium with the horizontal components of the force in the shell at $\mathrm{r}=\mathrm{a}$. Hence from Fig. 1 and Fig. 3, we have:
$N_{x} \operatorname{Cos} \phi_{o}+\bar{Q}_{x v} \operatorname{Sin} \phi_{o}=\bar{Q}_{x}$
Here $\bar{Q}_{X v}$ and $\bar{Q}_{x}$ are the equivalent transverse shear forces for shell and nozzle respectively, including the effects of the twisting moments. Since $\phi_{O}$ is assumed to be small, this condition reduces to
$\left(1-a^{2} / 2 R^{2}\right) N_{x}+(a / R) \bar{Q}_{x v}=\left(Q_{x}\right)_{x=0}$
Since boundary condition-III determine $C_{12}$, the rest six unknowns $\mathrm{H}_{1}, \mathrm{H}_{2}, \mathrm{H}_{3}, \mathrm{H}_{4}, \quad \mathrm{C}_{3}$ and $\mathrm{C}_{4}$ are to be found using conditions I, II and IV through VII.

From conditions I \& II葉, we set [2]:
$\mathrm{H}_{4}=-\mathrm{H}_{1}-(\mu \mathrm{M} / \pi \mathrm{ahE})$
and $\mathrm{H}_{3}=\mathrm{H}_{5} / \mathrm{a}$
Using the remaining 4 conditions (i.e. IV through VII), the constants $\mathrm{C}_{3}, \mathrm{C}_{4}, \mathrm{H}_{1}$ and $\mathrm{H}_{2}$ can be found from equations (54), (59), (61) and (68) of Ref. [2]. These equations are numbered (15) to (18) in this paper and are given below:

$$
\begin{align*}
& D_{1} C_{3}+D_{2} C_{4}-\left[\alpha_{1}-\left(L_{1} / a\right)\right] H_{1}-\left[\beta_{1}+\left(L_{2} / a\right)\right] H_{2}=0  \tag{15}\\
& D_{3} C_{3}+D_{4} C_{4}-\left[\left(1+L_{3}\right) \gamma\left(12\left(1-\mu^{2}\right)\right)^{\left.\frac{1}{2} / \rho u^{2}\right] H_{1}-\left[L_{4} \gamma\left(12\left(1-\mu^{2}\right)\right)^{1 / 2} /\right.}\right. \\
& \left.\rho u^{2}\right] H_{2}=(1 / \pi \mu)\left[\left(\mu \rho^{2} / \gamma\right)-\left((1+\mu) \cdot\left(12\left(1-\mu^{2}\right)\right)^{\left.\left.\frac{1}{2} / u^{2}\right)\right] \cdot\left(R M / \mathrm{Et}^{2} 1\right)}\right.\right. \tag{16}
\end{align*}
$$

$\mathrm{D}_{5} \mathrm{C}_{3}+\mathrm{D}_{6} \mathrm{C}_{4}-\mathrm{D}_{7} \mathrm{H}_{1}+\mathrm{D}_{8} \mathrm{H}_{2}=\left[\left(2-\mu^{2}\right) /\left(12\left(1-\mu^{2}\right) \pi\right)\right] \cdot\left(\mathrm{RM} / \mathrm{Et}^{2} 1\right)$
$\mathrm{D}_{9} \mathrm{C}_{3}+\mathrm{D}_{10} \mathrm{C}_{4}+\mathrm{D}_{11} \mathrm{H}_{1}-\mathrm{D}_{12} \mathrm{H}_{2}=\left[(1-\eta)\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}} / \pi \mathrm{u}^{3}\right]$. ( $\mathrm{RM} / \mathrm{Et}^{2} 1$ )
where constants $D_{1}$ to $D_{12}$ are given below:
$D_{1}=(u / a) \cdot\left[K e i u+(2 / u) \operatorname{Ker}^{\prime} u\right]$
$D_{2}=(u / a)\left[(2 / u) \cdot K e i^{\prime} u-K e r u\right]$
$D_{3}=-[(1+\mu) / u]$. Ker $u+$ Ker $^{\prime} u+\left[2\left(1+\mu / u^{2}\right] \cdot K e i^{\prime} u\right.$
$D_{4}=-[(1+\mu) / u]$. Kei $u+$ Kei $^{\prime} u-\left[2(1+\mu) u^{2}\right] \cdot$ Kezíu
$D_{5}=\left[\gamma^{2} u /\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}}\right] \cdot\left[(1-\mu) u^{-1} \cdot\right.$ Kei $u+$
$2(1-\mu) u^{-2} \cdot$ Ker $\left.^{\prime} u-K e i^{\prime} u\right]$
$D_{6}=\left[\gamma^{2} u /\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}}\right] \cdot\left[-(1-\mu) u^{-1} \cdot\right.$ Ker $u$
$+2(1-\mu) u^{-2} \cdot$ Kei $\left.^{\prime} u+K^{\prime} r^{-} u\right]$
$D_{7}=\left[\gamma^{2} a^{2} / \rho^{3} u\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}}\right] \cdot\left[\alpha_{1}^{2}-\beta_{1}^{2}-\mu\left(1+L_{3}\right) / a^{2}\right.$
$\left.+\left(\alpha_{1} L_{1}+\beta_{1} L_{2}\right) / a\right]$
$D_{8}=\left[\gamma^{2} a^{2} /\left(\rho^{3} u\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}}\right] \cdot\left[2 \alpha_{1} \beta_{1}+\right.\right.$
$\left.\mu L_{4} / a^{2}-\left(\alpha_{1} L_{2}-\beta_{1} L_{1}\right) / a\right]$
$D_{9}=(1-\eta) u^{-1}\left(\operatorname{Ker~} u-2 \mathrm{u}^{-1} \operatorname{Ke} 1^{\prime} \mathrm{u}\right)-$
$\left[u^{3} \rho^{2} /\left(12\left(1-\mu^{2}\right) \gamma^{2}\right)\right] \cdot\left[(1-\mu) u^{-2}\right.$ Kei u
$\left.-\operatorname{Ker} u+u^{-1} \cdot K e i^{\prime} u+2(1-\mu) \cdot u^{-3} \cdot \operatorname{Ker}^{\prime} u\right]$
$D_{10}=(1-\eta) \cdot u^{-1} \cdot\left(\right.$ Kei $u+2 u^{-1} \cdot$ Ker $\left.^{\prime} u\right)+$
$\left[u^{3} \rho^{2} /\left(12\left(1-\mu^{2}\right)\right) \gamma^{2}\right] \cdot\left[(1-\mu) u^{-2}\right.$. Ker $u$

$D_{11}=\left[a^{3} /\left(u^{2} \rho^{2} \gamma\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}}\right] \cdot\left[\alpha_{1}\left(\alpha_{1}^{2}-3 \beta_{1}^{2}\right)-\right.\right.$
$(2-\mu) \alpha_{1} / a^{2}+(3-\mu)\left(-\alpha_{1} L_{3}-\beta_{1} L_{4}\right) /\left(2 a^{2}\right)+$
$\left.\left.\left(\alpha_{1}^{2}-\beta_{1}^{2}\right) L_{1}+2 \alpha_{1} \beta_{1} L_{2}\right) / a+(1-\mu) L_{1} /\left(2 a^{3}\right)\right]$
$D_{12}=\left[a^{3} / u^{2} \rho^{2} \gamma\left(12\left(1-\mu^{2}\right)\right)^{\frac{1}{2}}\right] \cdot\left[\beta_{1}\left(3 \alpha^{2}-\beta_{1}^{2}\right)-\right.$
$(2-\mu) \beta_{1} / a^{2}-(3-\mu)\left(-\alpha_{1} L_{4}+\beta_{1} L_{3}\right) /\left(2 a^{2}\right)-$
$\left.\left(\left(\alpha_{1}^{2}-\beta_{1}^{2}\right) L_{2}-2 \alpha_{1} \beta_{1} L_{1}\right) / a-(1-\mu) L_{2} /\left(2 a^{3}\right)\right]$
where $\alpha_{1}$ and $\beta_{1}$ are given by Eq. (11) of Ref. [2] and are as follows:
$\alpha_{1}=(1 / a) \cdot\left[1-1 /(2 \mu)+\left(3\left(1-\mu^{2}\right) \gamma^{2}+1-3 /\left(4 \mu^{2}\right)\right)^{\frac{1}{2}}\right]^{\frac{1}{2}}$
$\beta_{1}=(1 / a) \cdot\left[-(1-(1 / 2 \mu))+\left(3\left(1-\mu^{2}\right) \gamma^{2}+1-3 /\left(4 \mu^{2}\right)\right)^{\frac{1}{2}}\right]^{\frac{1}{2}}$
The constants $\mathrm{L}_{1}$ to $\mathrm{L}_{4}$ are as follows:
$L_{1}=K_{1} / K_{5}$
$L_{2}=K_{2} / K_{5}$
$L_{3}=K_{3} / K_{5}$
$L_{4}=K_{4} / K_{5}$
where constants $\mathrm{K}_{1}$ to $\mathrm{K}_{5}$ are as given below:
$K_{1}=a \alpha_{1}\left[(1-\mu) \cdot\left(1-3 \mu\left(1+\mu+\mu^{2}\right)\right)-12\left(1-\mu^{4}\right) \gamma^{2}+\right.$
$\left.\left(2(1-\mu) \cdot\left(2+3 \mu+3 \mu^{2}\right)+24 \mu\left(1-\mu^{2}\right) \gamma^{2}\right) \mathrm{a}^{2} \beta_{1}^{2}\right]$
$K_{2}=-a \beta_{1}\left[(1-\mu) \cdot\left(1-3 \mu\left(1+\mu+\mu^{2}\right)\right)-12\left(1-\mu^{4}\right) \gamma^{2}-\right.$
$\left.\left(2(1-\mu) \cdot\left(2+3 \mu+3 \mu^{2}\right)+24 \mu\left(1-\mu^{2}\right) \gamma^{2}\right) a^{2} \alpha_{1}^{2}\right]$
$K_{3}=(1-\mu)(3+\mu)\left(1-3 \mu^{2}\right)+12\left(1-2 \mu-\mu^{2}\right)\left(1-\mu^{2}\right) \gamma^{2}$
$K_{4}=2\left[\left(4+9 \mu+3 \mu^{2}\right)(1-\mu)+12(2+\mu)\left(1-\mu^{2}\right) \gamma^{2}\right] a^{2} \alpha_{1} \beta_{1}$
$K_{5}=\left[12\left(1-\mu^{2}\right) \gamma^{2}+(1-\mu) \cdot(1+3 \mu)\right]^{2}$
Thus the final solution for the deflection of the spherical shell due to a bending moment M is given by equation (1) where constants $C_{3}$ and $C_{4}$ are obtained by solution of 4 simultaneous equations (15) to (18) for 4 unknowns; namely $\mathrm{C}_{3}, \mathrm{C}_{4}, \mathrm{H}_{1}$ and $\mathrm{H}_{2}$. The other constants; namely $\mathrm{D}_{1}$ to $\mathrm{D}_{12}, \mathrm{~L}_{1}$ to $\mathrm{L}_{4}, \mathrm{~K}_{1}$ to $\mathrm{K}_{5}$, $\alpha_{1}$ and $\beta_{1}$ used in these equations are given by equations (19) to (41).
approach for evaluation of rotational spring constants
The expression for deflection
$\mathrm{w}_{\mathrm{v}}=\left(\mathrm{C}_{3} \mathrm{Ker}^{\prime} \mathrm{s}+\mathrm{C}_{4} \mathrm{Kei}^{\prime} \mathrm{s}\right) \operatorname{Cos} \theta$
is valid at any point of spherical pressure vessel specified by radius $r$. Since, we are interested only in the deflection of the vessel at the juncture of nozzle and spherical shell, we let $\mathrm{r}=\mathrm{a}$.

At $r=a, s=a / 1=u$
Also, for present analysis $\theta=0$ since we are interested in maximum deflection due to moment which occurs in the plane where the moment is acting. Hence, the expression for $w_{v}$ for this case becomes:
$w_{v}=C_{3} \operatorname{Ker}^{\prime} u+C_{4} K^{K e} i^{\prime} u$
where $u=a / 1=a .\left[12\left(1-\mu^{2}\right) / R^{2} t^{2}\right]^{\frac{1}{4}}$
Substituting $\mu=0.3$, we get:
$\mathrm{u}=1.81784 . \mathrm{a} / \sqrt{\mathrm{Rt}}$ or
$u=1.81784 .(a / R) \sqrt{R / t}$
The rotational spring constant due to bending moment is given by referring to Fig. 4:
$K_{B}=M /|\phi|$
From Fig. 4:
$\mathrm{w}_{\mathrm{v}}=\phi \cdot \mathrm{a}$
$|\phi|=\left|w_{v}\right| / a$
Substituting this expression for $\phi$ in equation (44), we get:
$K_{B}=M \cdot a /\left|w_{v}\right|$


Fig. 4 Deflection (Wv) due to a bending moment (M)

Taking $M=1.0$ in $-1 b$, and substituting the expression for $w_{v}$ from equation (42), the expression for rotational spring constant becomes:
$\mathrm{K}_{\mathrm{B}}=\mathrm{a} /\left(\mathrm{C}_{3} \mathrm{Ker}^{\prime} \mathrm{u}+\mathrm{C}_{4} \mathrm{Kei}^{\prime} \mathrm{u}\right)$
For a given set of shell parameters ( $R \& t$ ) and nozzle parameters ( $a \& h$ ), unique values of $u, C_{3}$ and $C_{4}$ exist. Hence, given $R, t, a$ and $h$, unique values of deflection $\left(w_{V}\right)$ and rotational spring constant ( $K_{B}$ ) can be obtained. For convenience, we define 3 independent dimensionless parameters $\beta, \gamma$ and $\rho$ as follows:

$$
\begin{array}{ll}
\text { Shell Parameter: } & \beta=\mathrm{R} / \mathrm{t} \\
\text { Nozzle Parameter: } & \gamma=\mathrm{a} / \mathrm{h}
\end{array}
$$

Shell to nozzle thickness ratio: $\quad \rho=t / h$
Fixing the values of $\beta, \gamma$ and $\rho$ does not fix $R$, $t, a$ and $h$ e.g. If $R, t, a$ and $h$ are doubled, the ratios $\beta, \gamma$ and $\rho$ still remain the same while the rotational spring constant ( $K_{B}$ ) for this new set of values sets changed. Hence, to arrive at a unique value of $\mathrm{K}_{\mathrm{B}}$, we must fix one more parameter apart from $\beta, \gamma$ and $\rho$. This is achieved as follows:

For given values of $R, t, a$ and $h$, we find the value of $u$. Since, we are interested at the juncture of shell and nozzle, using $u=s$ curve [2], we read a constant value for $w_{v} E \cdot t^{2} /(M \sqrt{R / t} \cdot \operatorname{Cos} \theta)$.

$$
\text { Hence, } \frac{w_{v} E \cdot t^{2}}{M \sqrt{R / t} \cdot \cos \theta}=C \text { (Constant) }
$$

For our case $\theta=0$. Hence we have:

$$
\frac{w_{v} \mathrm{E} \cdot \mathrm{t}^{2}}{\mathrm{M} \sqrt{\mathrm{R} / \mathrm{t}}}=C
$$

$$
\begin{equation*}
\frac{M}{w_{v}}=\frac{E \cdot t^{2}}{c \cdot \sqrt{R / t}} \tag{47}
\end{equation*}
$$

From equation (45), we have the spring constant (K) as follows:

$$
\begin{align*}
& K_{B}=\frac{M \cdot a}{\left|w_{v}\right|}=\frac{E \cdot t^{2} \cdot a}{C \cdot \sqrt{R / t}} \\
& K_{B}=E \cdot t^{2} \cdot a /(C \cdot \sqrt{R / t}) \tag{48}
\end{align*}
$$

where $E$ is the Young's Modulus.
Equation (48) can also be written as:

$$
\begin{equation*}
K_{B} /\left(a \cdot t^{2}\right)=C^{\prime} / \sqrt{B} \tag{49}
\end{equation*}
$$

where $\beta=R / t$ (as per definition)
and $C^{\prime}=E / C=$ Another Constant
Hence, to fix the value of $K_{B}$ for fixed values of $B, \gamma$ and $\rho$, we must fix a and $t$ also.

It logically follows from equation (49) that there exists a unique value of $K_{B} /\left(a \cdot t^{2}\right)$ for a given set of values $\beta, \gamma$ and $\rho$.

At this stage, it seems logical to plot the values of $K_{B} / a \cdot t^{2}$ vs. $B$ for various combinations of $\gamma$ and $\rho$ because there exists a unique value of $K_{B} /\left(a \cdot t^{2}\right)$ for a given set of $\beta, \gamma$, and $\rho$. In fact, in latter part of this paper, values for $K_{B} /\left(a . t^{2}\right)$ have been plotted against the shell parameter $\beta(R / t)$ for 3 different values of the nozzle parameter $\gamma(\mathrm{a} / \mathrm{h})$ and unique values of the constant $\rho(t / h)$.

Using all the above equations, a computer program has been written to evaluate spring constants ( $K_{B}$ ) and ratios $K_{B} /\left(a . t^{2}\right)$ for various combinations of $R, t, a \& h$ (i.e. $\beta, \gamma$ and $\rho$ ). The various equations for Kelvin functions (ker $u, \operatorname{Kei} u$ ) and their derivatives (Ker ${ }^{\prime} u$, Kei'u) are taken from Ref. [3].

Computer results for spring constants $K_{B}$ and ratio $K_{B} /\left(a \cdot t^{2}\right)$ are obtained for various values of $\beta, \gamma$ and $\rho$. Tables 1 through 4 give ranges for $\rho=1.0,2.0,4.0,10.0$; $\gamma=5,10,15$ and various ranges of $\beta$.

The values of deflection ( $w_{V}$ ) computed in this study have been verified with Bijlaard's work in Ref. [2] and are found to be matching quite closely. Values taken from Tables 1 through 4 are also plotted in this paper to facilitate the vessel design and stress analysis. These are given in Fig. 5 to 8 respectively.

One must note that the values of rotational spring constants are valid only if the deflections are limited to a segment of shell that can be considered shallow. This leads to conditions:
(1) $u \leq 1.0, \mathrm{R} / \mathrm{t} \geq 10$
(2) $\mathrm{u}>1.0, \mathrm{R} / \mathrm{t} \geq 10$

Also the values of spring constants are considered accurate only of $a / R$ is less than or equal to $1 / 3$. This leads to the condition:
(3) $a / R \leq 1 / 3$, i.e $\mu \leq 1.05 \sqrt{\gamma / \rho}$

The above three conditions are derived in Ref. [4] and can be found in Appendix I. The computer program is designed to take care of these 3 conditions.

## CORRECTION OF THE SPRING CONSTANTS FOR HOT MODULUS

As we know the value of $E$ varies depnding upon the temperature of operation. The graphs plotted here are also for cold condition i.e. $E=30 \mathrm{E} 06 \mathrm{psi}$.

In case, the temperature of operation is higher than the normal ambient temperature, the value of bending spring constant ( $K_{B}$ as obtained at ambient temperature) must be corrected as follows:

As given by equation (48), we know:
$K_{B}=E \cdot t^{2} \cdot a /(C \cdot \sqrt{R / t})$
For a given geometry of spherical shell:
$K_{B} \propto E$
$\frac{K_{B} \text { (at temperature of operation) }}{K_{B} \text { (at ambient temperature) }}=\frac{E_{h}}{E_{C}}$
or $\quad K_{B}$ (at temperature of operation $=$
$\left(E_{h} / E_{c}\right) \cdot K_{B}$ (at ambient temperature)
$\mathrm{E}_{\mathrm{c}}=30 \mathrm{E} 06$
$E_{h}=$ Modulus of Elasticity at temperature of operation

Thus, if $\beta, \gamma, \rho, t, a$ and temperature of operation are specified, the spring constant can readily be found using these curves. Given below is an example to illustrate the use of these curves:

## NUMERICAL EXAMPLE

## Given:

$R=500^{\prime \prime}, t=2^{\prime \prime}, a=5^{\prime \prime}, h=1^{\prime \prime}$
Temperature of operation $=600 \mathrm{~F}$
Material : C-Steel with Carbon Content $\leq 0.3 \%$

## Required:

(a) Rotational spring constant $\left(K_{B}\right)$ at room temperature
(b) Rotational spring constant ( $\mathrm{K}_{\mathrm{B}}$ ) at 600 F

## Solution:

$\beta=(R / t)=500 / 2=250$
$\gamma=(a / h)=5 / 1=5$
$\rho=(t / h)=2 / 1=2$
We refer to Fig. 6 corresponding to the value of $\rho=2$. On this graph, we select the curve with $\gamma=5$. For $\beta=250$, we read from the curve:

$$
K_{B} /\left(\mathrm{a} \cdot \mathrm{t}^{2}\right)=0.675 \mathrm{E} 07 \mathrm{lbs} / \mathrm{in}^{2}
$$

Substituting $a=5^{\prime \prime}$ and $t=2^{\prime \prime}$, we set:

$$
K_{B}=0.675 \mathrm{E} 07 \times 5 \times 2^{2}=13.5 \mathrm{E} 07 \mathrm{in}-1 \mathrm{~b} / \mathrm{rad} .
$$

This value of rotational spring constant ( $K_{B}$ ) is at ambient temperature (say 70 F ). The bending spring constant ( $\mathrm{K}_{\mathrm{B}}$ ) at 600 F is obtained as follows:

$$
\text { At } \begin{aligned}
\mathrm{T} & =70 \mathrm{~F}, \mathrm{E}_{\mathrm{c}}=30 \mathrm{E} 06 \mathrm{psi} \\
\mathrm{~T} & =600 \mathrm{~F}, \mathrm{E}_{\mathrm{h}}=25.7 \mathrm{E} 06 \mathrm{psi}
\end{aligned}
$$

$K_{B}($ at 600 F$)=\left(\mathrm{E}_{\mathrm{h}} / \mathrm{E}_{\mathrm{c}}\right) \cdot \mathrm{K}_{\mathrm{B}}$ (at ambient temperature)

$$
=(25.7 / 30.0) \times 13.5 \mathrm{E} 07
$$

$K_{B}($ at 600 F$)=11.57 \mathrm{E} 07 \mathrm{in}-1 \mathrm{~b} / \mathrm{rad}$.

CONCLUSION

It has been found that rotational spring constant ( $\mathrm{K}_{\mathrm{B}}$ ) is inversely proportional to square root of the radius of the spherical shell and directional proportional to $t^{5 / 2}$ (where $t$ is the thickness of the spherical shell) and directly proportional to nozzle radius (a).

It can be seen from Tables 1 through 4 that as nozzle thickness (h) decreases, the bending spring constant ( $K_{B}$ ) decreases.

In order to limit the stresses at the juncture of radial nozzle and spherical shell, the design requirement is to reduce the rotational spring constant ( $\mathrm{KB}_{\mathrm{B}}$ ). Hence, to achieve a low value of rotational spring constant ( $K_{B}$ ), it is recommended:
(1) To increase the spherical shell radius ( $R$ )
(2) To reduce the spherical shell thickness ( $t$ )
(3) To reduce the nozzle radius (a)
(4) To reduce the nozzle thickness (h)

## REFERENCES

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5. Reissner, E., "Stresses and small displacements of shallow spherical shells". International Joural of Mathematical Physics, 25, pp. 80-85 (1946).
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## APPENDIX

DERIVATION OF CONDITIONS FOR VALIDITY OF GRAPHS:
To check the validity of graphs, we make the following 2 derivations:

Derivation I:
As suggested by Bijlaard [6], the graphs are applicable only if deflections are limited to a segment of shell that can be considered as shallow, which still can be assumed if these deflections die out at a distance $r$ of about 0.6 R from the center of the attachment. The following two values of $s$ are also as per Ref. [6]:
(a) For $u(a / 1) \leq 1.0$, Deflections die out at about $s(\bar{r} / 1)=3.5$
(b) For $u(a / 1)>1.0$, Deflections die out at about $s(r / 1)=u+2.5$

Since $s=1.81784 \cdot(r / R) \cdot \sqrt{R / t}$, we get

$$
\begin{equation*}
r=0.55 \sqrt{R}{ }^{R} t . s \tag{50}
\end{equation*}
$$

Case (a)
$\mathrm{u} \leq 1.0, \mathrm{~s}=3.5$ (For deflections to die out)

From Eq. (50), we get:

$$
\begin{align*}
& \mathrm{r}=0.55 \sqrt{\mathrm{R} 7 \mathrm{t}} \times 3.5 \\
& \mathrm{r}=1.92 \sqrt{\mathrm{R} \frac{7}{} \mathrm{t}} \tag{51}
\end{align*}
$$

Since $r \leq 0.6 R$ (For deflection to die out), we set from Eq. (51) :

$$
1.92 \sqrt{\mathrm{R} 7 \mathrm{t}} \leq 0.6 \mathrm{R}
$$

$$
\begin{equation*}
\sqrt{\mathrm{R} / \mathrm{t}} \geq 1.92 / 0.6=3.21 \tag{52}
\end{equation*}
$$

or $\quad R / t \geq 10$
Case (b)
$u>1.0, s=u+2.5$ (For deflections to die out)
From Eq. (50), we get

$$
\begin{equation*}
r=0.55 \sqrt{R \neq t} .(u+2.5) \tag{53}
\end{equation*}
$$

Since $r \leq 0.6 R$ (For deflections to die out), we get from Eq. (53) :

$$
0.55 \sqrt{R / t} \cdot(u+2.5) \leq 0.6 R
$$

or $\sqrt{R / t} \geq(0.9167 u+2.292)$
By approximation, we set:

$$
\begin{equation*}
\mathrm{R} / \mathrm{t} \geq(u+2.3)^{2} \tag{54}
\end{equation*}
$$

Derivation II:
As suggested by Bijlaard [6], the graphs may only be expected to be sufficiently accurate if $a / R$ is less than or equal to $1 / 3$.

From Eq. (20) of Ref. [ㅢㅡ]:

$$
\begin{equation*}
a / R=\left(u^{2} \cdot \rho / \gamma\right) /\left[12\left(1-\mu^{2}\right)\right]^{\frac{1}{2}} \tag{55}
\end{equation*}
$$

(Note: It can be verified by substituting the values of $u, p \& \gamma$ in right side of above equation that it is equal to $a / R$ )

Hence from Eq. (55), we get:

$$
\left(u^{2} \cdot \rho / \gamma\right) /\left[12\left(1-\mu^{2}\right)\right]^{\frac{1}{2}} \leq 1 / 3
$$

Substituting $\mu=0.3$, we get:

$$
\begin{align*}
& \mathrm{u}^{2} \leq 1.101 \cdot(\gamma / \rho) \\
& \mathrm{u} \leq 1.05 \sqrt{\gamma / \rho} \tag{56}
\end{align*}
$$



Fig. 5 Graph for rotational spring constants for $p(t / h)=1.0$


Fig. 6 Graph for rotational apring constants for $\rho(\mathrm{t} / \mathrm{h})=2.0$


Fig. 7 Graph for rotational spring constants for $\rho(\mathrm{t} / \mathrm{h})=4.0$


Fig. 8 Graph for rotational spring constants for $\rho(t / h)=10.0$

TABLE 1: RESULTS FOR ROTATIONAL SPRING CONSTANTS FOR RHO $=1.0$

| $\mathrm{K}_{\mathrm{B}}$-Spring constant |  |  |  | $\begin{gathered} \text { BETA }=\mathrm{R} / \mathrm{t} \\ \hline \text { DEFL. } \end{gathered}$ | GAMMA $=a / h$$K_{B}$ | RHO $=\mathrm{t} / \mathrm{h} \quad u=1.81784 * \mathrm{a}$ *SQRT (BETA) $/ \mathrm{R}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| R | t | a | h |  |  | $\mathrm{K}_{\mathrm{B}} /(\mathrm{a} * \mathrm{t} * * 2)$ | BETA | gamma | u |
| 150. | 1.0 | 5.0 | 1.0 | -0.000000098 | 0.50934 E 08 | 0.10187 E 08 | 150. | 5.0 | 0.742130 |
| 320. | 1.6 | 8.0 | 1.6 | -0.000000045 | 0.17741 E 09 | 0.86628 E 07 | 200. | 5.0 | 0.642703 |
| 500. | 2.0 | 10.0 | 2.0 | -0.000000032 | 0.30897 E 09 | 0.77243 E 07 | 250. | 5.0 | 0.574851 |
| 720. | 2.4 | 12.0 | 2.4 | -0.000000025 | 0.48931 E 09 | 0.70792 E 07 | 300. | 5.0 | 0.524764 |
| 1050. | 3.0 | 15.0 | 3.0 | -0.000000017 | 0.89148 E 09 | 0.66035 E 07 | 350. | 5.0 | 0.485838 |
| 1440. | 3.6 | 18.0 | 3.6 | -0.000000012 | 0.14546 E 10 | 0.62353 E 07 | 400. | 5.0 | 0.454459 |
| 1800. | 4.0 | 20.0 | 4.0 | -0.000000011 | 0.19008 E 10 | 0.59399 E 07 | 450. | 5.0 | 0.428469 |
| 2200. | 4.4 | 22.0 | 4.4 | -0.000000009 | 0.24262 E 10 | 0.56964 E 07 | 500. | 5.0 | 0.406481 |
| 150. | 1.0 | 10.0 | 1.0 | -0.000000066 | 0.15218 E 09 | - 0.15218 E 08 | 150. | 10.0 | 1.484259 |
| 320. | 1.6 | 16.0 | 1.6 | -0.000000033 | 0.48058 E 09 | 0.11733 E 08 | 200. | 10.0 | 1.285406 |
| 500. | 2.0 | 20.0 | 2.0 | -0.000000026 | 0.77936 E 09 | 0.97419 E 07 | 250. | 10.0 | 1.149702 |
| 720. | 2.4 | 24.0 | 2.4 | -0.000000021 | 0.11683 E 10 | 0.84512 E 07 | 300. | 10.0 | 1.049528 |
| 1050. | 3.0 | 30.0 | 3.0 | -0.000000015 | 0.20369 E 10 | 0.75440 E 07 | 350. | 10.0 | 0.971676 |
| 1440. | 3.6 | 36.0 | 3.6 | -0.000000011 | 0.32050 E 10 | 0.68695 E 07 | 400. | 10.0 | 0.908919 |
| 1800. | 4.0 | 40.0 | 4.0 | -0.000000010 | 0.40620 E 10 | 0.63470 E 07 | 450. | 10.0 | 0.856937 |
| 2200. | 4.4 | 44.0 | 4.4 | -0.000000009 | 0.50507 E 10 | 0.59292 E 07 | 500. | 10.0 | 0.812962 |
| 150. | 1.0 | 15.0 | 1.0 | -0.000000042 | 0.35928 E 09 | 0.23952 E 08 | 150. | 15.0 | 2.226388 |
| 320. | 1.6 | 24.0 | 1.6 | -0.000000022 | 0.10670 E 10 | 0.17367 E 08 | 200. | 15.0 | 1.928108 |
| 500. | 2.0 | 30.0 | 2.0 | -0.000000018 | 0.16501 E 10 | 0.13751 E 08 | 250. | 15.0 | 1.724553 |
| 720. | 2.4 | 36.0 | 2.4 | -0.000000015 | 0.23805 E 10 | 0.11480 E 08 | 300. | 15.0 | 1.574293 |
| 1050. | 3.0 | 45.0 | 3.0 | -0.000000011 | 0.40197 E 10 | 0.99253 E 07 | 350. | 15.0 | 1.457513 |
| 1440. | 3.6 | 54.0 | 3.6 | -0.000000009 | 0.61554 E 10 | 0.87955 E 07 | 400. | 15.0 | 1.363378 |
| 1800. | 4.0 | 60.0 | 4.0 | -0.000000008 | 0.76200 E 10 | 0.79375 E 07 | 450. | 15.0 | 1.285405 |
| 2200. | 4.4 | 66.0 | 4.4 | -0.000000007 | 0.92812 E 10 | 0.72636 E 07 | 500. | 15.0 | 1.219442 |

$$
\begin{array}{lll}
\text { UNITS : } & \text { R(IN.) } & \text { a (IN.) }
\end{array} \quad \text { DEFL. (IN.) } \quad \text { (IN.) } \quad \text { h (IN.) } \quad 1 \text { K (IN.-LB/RAD.) }
$$

TABLE 2: RESULTS FOR RATATIONAL SPRING CONSTANTS FOR RHO $=2.0$

| R | t | a | h | DEFL. | $\mathrm{K}_{\mathrm{B}}$ | $\mathrm{K}_{\mathrm{B}} /(\mathrm{a} * \mathrm{t} * * 2)$. | BETA | GAMMA | u |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 200. | 2.0 | 5.0 | 1.0 | -0.000000028 | 0.17716 E 09 | $0.88579 \mathrm{E} \quad 07$ | 100. | 5.0 | 0.454460 |
| 480. | 3.2 | 8.0 | 1.6 | -0.000000013 | 0.63614 E 09 | $0.77654 \mathrm{E} \quad 07$ | 150. | 5.0 | 0.371065 |
| 800. | 4.0 | 10.0 | 2.0 | -0.000000009 | $0.11436 \mathrm{E} \quad 10$ | 0.71476 E 07 | 200. | 5.0 | 0.321351 |
| 1080. | 4.8 | 12.0 | 2.4 | -0.000000006 | $0.19143 \mathrm{E} \quad 10$ | $0.69237 E 07$ | 225. | 5.0 | 0.302973 |
| 1500. | 6.0 | 15.0 | 3.0 | -0.000000004 | $0.36372 \mathrm{E} \quad 10$ | 0.67356 E 07 | 250. | 5.0 | 0.287426 |
| 2160. | 7.2 | 18.0 | 3.6 | -0.000000003 | 0.60044 E 10 | 0.64347 E 07 | 300. | 5.0 | 0.262382 |
| 2800. | 8.0 | 20.0 | 4.0 | -0.000000003 | 0.79382 E 10 | 0.62017 E 07 | 350. | 5.0 | 0.242919 |
| 3520. | 8.8 | 22.0 | 4.4 | -0.000000002 | $0.10246 \mathrm{E} \quad 11$ | 0.60139 E 07 | 400. | 5.0 | 0.227230 |
| 200. | 2.0 | 10.0 | 1.0 | -0.000000027 | 0.36402 E 09 | $0.91006 \mathrm{E} \quad 07$ | 100. | 10.0 | 0.908920 |
| 480. | 3.2 | 16.0 | 1.6 | -0.000000014 | 0.11825 EE 10 | 0.72173 E 07 | 150. | 10.0 | 0.742130 |
| 800. | 4.0 | 20.0 | 2.0 | -0.000000010 | 0.20034 E 10 | 0.62606 E 07 | 200. | 10.0 | 0.642703 |
| 1080. | 4.8 | 24.0 | 2.4 | -0.000000007 | 0.32813 E 10 | 0.59340 E 07 | 225. | 10.0 | 0.605946 |
| 1500. | 6.0 | 30.0 | 3.0 | -0.000000005 | $0.61216 \mathrm{E} \quad 10$ | 0.56681 E 07 | 250. | 10.0 | 0.574851 |
| 2160. | 7.2 | 36.0 | 3.6 | -0.000000004 | $0.98136 \mathrm{E} \quad 10$ | 0.52585 E 07 | 300. | 10.0 | 0.524765 |
| 2800. | 8.0 | 40.0 | 4.0 | -0.000000003 | 0.12684 E 11 | $0.49547 \mathrm{E} \quad 07$ | 350. | 10.0 | 0.485838 |
| 3520. | 8.8 | 44.0 | 4.4 | -0.000000003 | 0.16077 E 11 | 0.47184 E 07 | 400. | 10.0 | 0.454459 |
| 200. | 2.0 | 15.0 | 1.0 | -0.000000022 | 0.68538 E 09 | 0.11423 E 08 | 100. | 15.0 | 1.363379 |
| 480. | 3.2 | 24.0 | 1.6 | -0.000000012 | 0.20405 E 10 | 0.83029 E 07 | 150. | 15.0 | 1.113194 |
| 800. | 4.0 | 30.0 | 2.0 | -0.000000009 | 0.32712 E 10 | $0.68150 \mathrm{E} \quad 07$ | 200. | 15.0 | 0.964055 |
| 1080. | 4.8 | 36.0 | 2.4 | -0.000000007 | 0.52468 E 10 | 0.63257 E 07 | 225. | 15.0 | 0.908919 |
| 1500. | 6.0 | 45.0 | 3.0 | -0.000000005 | 0.96150 E 10 | 0.59352 E 07 | 250. | 15.0 | 0.862277 |
| 2160. | 7.2 | 54.0 | 3.6 | -0.000000004 | 0.14973 E 11 | 0.53487 E 07 | 300. | 15.0 | 0.787147 |
| 2800. | 8.0 | 60.0 | 4.0 | -0.000000003 | 0.18918 E 11 | 0.49264 E 07 | 350. | 15.0 | 0.728757 |
| 3520. | 8.8 | 66.0 | 4.4 | -0.000000003 | 0.23541 E 11 | 0.46060 E 07 | 400. | 15.0 | 0.681689 |
| UNITS : |  | $\begin{aligned} & \text { R (IN.) } \\ & \mathrm{t}(\mathrm{IN} .) \end{aligned}$ |  | $\begin{aligned} & a(I N .) \\ & h \text { (IN.) } \end{aligned}$ | $\begin{aligned} & \text { DEFL. } \\ & K_{B}(I N . \end{aligned}$ |  |  |  |  |

TABLE 3: RESULTS FOR ROTATIONAL SPRING CONSTANTS FOR RHO $=4.0$

table 4: Results for rotational spring constants for rho $=10.0$


